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Electron-nucleus interactions and nuclear effects in atomic transitions

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EJC Oleron, October 2021





NATURWISSENSCHAFTLICHE FAKULTÄT



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"Nuclear"	effects			



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- small corrections to atomic level and transition energies
- best studied by comparing two different isotopes
 → measuring isotope shifts

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Bridging atomic and nuclear physics



- exploring nuclear properties via atomic physics experiments
- nuclear processes directly involving atomic electrons

AP, Contemporary Physics 51, 471 (2010)

The borderline between atomic and nuclear physics

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Bridging atomic and nuclear physics



- exploring nuclear properties via atomic physics experiments TODAY
- nuclear processes directly involving atomic electrons TOMORROW

AP, Contemporary Physics 51, 471 (2010)

The borderline between atomic and nuclear physics



nuclear transitions involving electrons





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Isotope shifts







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Isotope shifts

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Isotope	shifts			

... frequency difference in an electronic transition between two isotopes

$$\Delta \nu_{IS} = \Delta \nu_{FS} + \Delta \nu_{MS}$$





Fermi two-parameter charge distribution

$$\rho_{\textit{nuc}} = \frac{\rho_0}{1 + e^{(r-c)/a}}$$





- $\Delta \nu \sim Z^{5..6}$ • heavy nuclei $\Delta \nu / \nu = 10^{-5}$
- light nuclei $\Delta \nu / \nu = 10^{-8}$

numerical integration of the Dirac equation with V_{nuc}



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Relativistic nuclear recoil operator

$$\mathbf{R}_{ij} = \frac{\vec{p}_i \cdot \vec{p}_j}{2M} - \frac{Z\alpha}{2Mr_i} \left(\vec{\alpha}_i + \frac{(\vec{\alpha}_i \cdot \vec{r}_i)\vec{r}_i}{r_i^2}\right) \cdot \vec{p}_j$$

- normal mass shift correction $\langle \sum_i R_{ii} \rangle$
- specific mass shift term $\langle \sum_{i \neq j} R_{ij} \rangle$

in muonic atoms QED theory B. Fricke, PRL 30, 119 (1973)
V. M. Shabaev, Theor. Mat. Fiz. 63, 394 (1985)
Yad. Fiz. 47, 107 (1988)



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How to measure isotope shifts?



Talks by Ruben and Iain!

- muonic atoms
- electron scattering
- x-ray spectroscopy
- laser spectroscopy
- dielectronic recombination



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Highly charged ions



... storage ring ...

or EBIT



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Isotopic shifts in x-ray transitions



Figure 3. Spectra of the $2p_{3/2} \rightarrow 2s_{1/2}$ transitions in the highly charged ²³³U, ²³³U, and ²³⁸U isotopes. Individual lines are labeled by the charge state of the emitting ions: $C = U^{8i+}$, $B = U^{87+}$, $Be = U^{88+}$ and $Li = U^{80+}$.

$$\delta \langle r^2 \rangle^{233,238} = -0.457 \pm 0.043 \, \text{fm}^2$$



Elliot, Beiersdorfer, Chen, PRL 76, 1031 (1996)



Relativistic recoil and isotopic shifts in ${}^{40}Ar/{}^{36}Ar$



$$\begin{array}{ll} Ar^{13+} & 1s^2 2s^2 2p\,^2 P_{1/2}\, -^2\, P_{3/2} \\ Ar^{14+} & 1s^2 2s 2p\,^3 P_1\, -^3\, P_2 \end{array}$$

FIG. 1 (color online). A typical spectral line obtained from the $1s^22s^22p\ ^2P_{1/2}\ ^2P_{3/2}$ transition in B-like 40 Ar¹³⁺. The six dashed curves represent a fit to the Zeeman components.



sub-ppm accuracy: Soria Orts et al. PRL 97, 103002 (2006)

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Isotopic shifts ⁴⁰Ar/³⁶Ar

This experiment confirms the newest treatment of relativistic recoil effect

Ion	Theory 1	N _{Observed}	Isotopic shift	rs (⁴⁰ Ar/ ³⁶ Ar)
	(nm, air)*	(nm)	theory (nm)	experiment (nm)
Ar ¹³⁺	441.16(27)	441.2556(1)	0.00123(5)	0.00123(6)
Ar ¹⁴⁺	594.24(30)	594.3879(2)	0.00122(5)	0.00120(10)

Normal and specific mass shifts and their relativistic corrections are of similar sizes.

This relativistic few-body quantum problem can only be solved consistently within a full QED treatment.

courtesy J. Crespo



Soria Orts et al. PRL 97, 103002 (2006)

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Isotope shifts and nuclear halos





courtesy K. Blaum

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Isotope shifts and nuclear halos







Isotope shifts and nuclear halos





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\mathbf{RR}



- direct process
- any electron energy
- electron-radiation field



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- direct process
- any electron energy
- electron-radiation field





- resonant process
- Coulomb interaction
- Breit interaction



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Isotope shifts in DR



C. Brandau, C. Kozhuharov, Z. Harman et al., PRL 100, 073201 (2008)



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Isotope shifts in DR



C. Brandau, C. Kozhuharov, Z. Harman *et al.*, PRL **100**, 073201 (2008) nuclear deformation: Kozhedub *et al.*, PRA **77**, 032501 (2008)



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HFS





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Splittings or shifts of fine structure levels due to the interaction of nuclear multipole moments with the electromagnetic field created by the electrons at the nucleus

magnetic dipole moment associated to spin

$$\vec{F} = \vec{I} + \vec{J}$$



 electric quadrupole moment - deviation from spherical charge distribution



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Laser spectroscopy

Hyperfine splitting for some heavy H-like ions



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DR very o	close to thresh	old		

- HFS of 4s_{1/2} and 4p_{1/2} in ²⁰⁷Pb⁵³⁺ (comparing to ²⁰⁸Pb⁵³⁺) very low-energy electron captured in Rydberg state!
 R. Schuch, E. Lindroth *et al.*, PRL **95**, 183003 (2005)
- HFS of 2s state in ⁴⁵Sc¹⁸⁺ using DR Rydberg resonances M. Lestinsky, E. Lindroth *et al.*, PRL **95**, 183003 (2005)

TRICKS: low-energy electron and Rydberg state!

• hyperfine induced transitions: $2s2p \ {}^{3}P_{0} \rightarrow 2s^{2} \ {}^{1}S_{0}$ in Be-like ${}^{47}Ti^{18+}$

S. Schippers et al., PRL 98, 033001 (2007)



Nuclear hyperfine mixing in ²²⁹Th

The lowest known excited nuclear state at only 8 eV in ²²⁹Th



In ²²⁹Th⁸⁹⁺ the very strong 28 MT magnetic field of the unpaired electron mixes F = 2 states

V. M. Shabaev et al., arXiv:2109.01642 [physics.atom-ph]

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Nuclear polarization

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What does	s this mean?			

Due to exchange of virtual photons, nucleus undergoes virtual transitions to excited states!



$$\Delta E \sim (E_n - E_a)^{-1}$$

Main theoretical challenge for high-precision tests of QED!

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Nuclear level schemes





 $^{208}_{\ 82} \mathrm{Pb}_{126}$

200 meV for K-shell electron

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Nuclear level schemes





200 meV for K-shell electron

three orders of magnitude smaller!



































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